# The Structure of $\mathbf{S r P b}_{\mathbf{2}} \mathrm{I}_{\mathbf{6}} \cdot \mathbf{7} \mathbf{H}_{\mathbf{2}} \mathbf{O}$ 

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The structure of the salt $\mathrm{SrPb}_{2} \mathrm{I}_{6} \cdot 7 \mathrm{H}_{2} \mathrm{O}$ has been investigated. The compound crystallizes in the tetragonal system, with cell constants

$$
a=11 \cdot 77 \pm 0.02, c=17.32 \pm 0.02 \AA .
$$

The space group is $D_{4 h}^{19}-I 4_{1} / a m d$.
The structure is characterized by a three-dimensional framework of octahedra $\left[\mathrm{PbI}_{6 / 2}\right]$ joined to one another through all the corners. The empty spaces of this framework are occupied by the strontium ions and by the water molecules. Six water molecules octahedrally surround each strontium ion.

The errors in the intensities due to the particularly high absorbing power of the substance ( $\mu=962 \mathrm{~cm} .^{-1}$ ) do not affect the coordinates of the heavy atoms but do not allow an exact determination of the coordinates of the water molecules.

## Introduction

Strontium and lead(II) salts are often isotypic and form solid solutions. But in some compounds, such as $\mathrm{SrPb}_{2} \mathrm{I}_{6} .7 \mathrm{H}_{2} \mathrm{O}$ the two ions are combined in a definite ratio. The study of the salt $\mathrm{SrPb}_{2} \mathrm{I}_{6} .7 \mathrm{H}_{2} \mathrm{O}$ has been undertaken in order to detect any specific characteristics of the two elements. The research has shown that lead(II) forms partially covalent bonds with iodine and that strontium is surrounded by six water molecules, forming an octahedron of the type observed in simple hexahydrated salts.

## Experimental

Crystals were prepared by evaporating an aqueous solution of strontium iodide in the presence of undissolved lead iodide. Their habit is pseudooctahedral and their color ranges from orange to brown. The chemical analysis was in good agreement with those of Mosnier (1897) and consistent with the formula $\mathrm{SrPb}_{2} \mathrm{I}_{6} .7 \mathrm{H}_{2} \mathrm{O}$.

All the crystals used were small bipyramids. Photographs were taken with $\mathrm{Cu} K \alpha$ radiation ( $\bar{\lambda}=1 \cdot 5418 \AA$ ). Cell dimensions were determined from oscillation- and Weissenberg photographs (rotation axis [001]). The unit cell is tetragonal:

$$
a=11 \cdot 77 \pm 0.02, c=17.32 \pm 0.02 \AA
$$

and contains four stoichiometric units $\mathrm{SrPb}_{2} \mathrm{I}_{6} .7 \mathrm{H}_{2} \mathrm{O}$. The observed density $d_{4}^{26}=4 \cdot 32$ g.cm. ${ }^{-3}$ is somewhat lower than the calculated one $d_{c}=4.56$ g.cm..$^{-3}$.

Observed reflections were $h k l$ for $h+k+l=2 n$; $h k 0$ for $h=2 n$ and $k=2 n$; $h h l$ for $l=2 n$, which indicated the space group $D_{4 h}^{19}-I 4_{1} / a m d$.

Intensity data for the $h k 0$ and $0 k l$ zones were
evaluated photometrically from the Weissenberg pho. tographs using the multiple-film technique.

The absorption corrections ( $\mu=962 \mathrm{~cm} .^{-1}$ ) using the Henshaw method (1958) were limited to an average section for the $h k 0$ zone and to two sections for the $0 k l$ zone. The Lorentz and polarization factors were obtained graphically. The atomic scattering factors used were taken from Sagel (1958).

## Determination of the structure

Atomic coordinates were found by trial and error, most of the atoms being in special positions. Good agreement between calculated and observed intensities was found assuming six iodine atoms to be disposed octahedrally around each lead atom. Hypotheses based


Fig. 1. $\varrho_{0}(x, y)$. Contours are drawn at approximately $20 \mathrm{e} . \AA^{-2}$, starting with $20 \mathrm{e} . \AA^{-2}$.


Fig. 2. $\varrho_{0}(y, z)$. Contours are drawn at approximately 20 e. $\AA^{-2}$, starting with 20 e. $\AA^{-2}$.

Table 1. Observed and calculated structure factors for the hk0 reflections

| $h k 0$ | $F_{o}$ | $F_{c}$ |  | $h k 0$ | $F_{o}$ |
| ---: | ---: | ---: | ---: | ---: | ---: |
| 020 | - | -55 |  | 440 | 908 |
| 040 | 1240 | 1189 |  | 460 | 372 |
| 060 | 766 | -759 |  | -389 |  |
| 080 | 354 | 241 | $4,10,0$ | 554 | 532 |
| $0,10,0$ | 461 | -513 |  | $4,12,0$ | 254 |
| $0,12,0$ | 284 | 283 |  | $4,14,0$ | - |
| $0,14,0$ | - | -50 | $\vdots$ | 660 | - |
| 220 | 271 | -294 | 680 | -102 |  |
| 240 | 334 | 432 | $6,10,0$ | - | 35 |
| 260 | 286 | 276 | $6,12,0$ | 84 | -108 |
| 280 | 533 | 563 | $6,14,0$ | - | 91 |
| $2,10,0$ | 224 | 212 | 880 | 601 | 619 |
| $2,12,0$ | 227 | 182 | $8,10,0$ | - | 18 |
| $2,14,0$ | - | -35 | $8,12,0$ | 228 | 253 |
|  |  |  | $10,10,0$ | - | -101 |

The structure factors are calculated with the axis origin at $2 / \mathrm{m}$.
on a linear arrangement $\mathrm{I}-\mathrm{Pb}-\mathrm{I}$ were excluded. Six water molecules were assumed to surround octahedrally each strontium ion; the seventh water molecule can occupy the $(b)$ position only, which is at the centre of an empty space between the octahedra formed by iodine and lead.

The initial coordinates were confirmed and refined

Table 2. Observed and calculated structure factors for the 0kl reflections

| $0 k l$ | $F_{0}$ | $F_{c}$ | 0 kl | $F_{o}$ | $F_{\boldsymbol{c}}$ | 0kl | $F_{0}$ | $F_{c}$ | 0 kl | $F_{0}$ | $F_{c}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 004 | 295 | -413 | 0,3,15 | - | 38 | 0,6,18 | - | 104 | 0,10,6 | 275 | 186 |
| 008 | 773 | 986 | 0,3,17 |  | -50 | 0,6,20 | - | 96 | 0,10,8 | 85 | -129 |
| 0,0,12 | - | - 145 | 0,3,19 | 162 | -244 | 071 | - | 16 | 0,10,10 | 234 | 146 |
| 0,0,16 | 440 | 586 | 0,3,21 | - | 85 | 073 | 477 | -486 | 0,10,12 | 140 | 146 |
| 0,0,20 |  | -78 | 040 | 948 | 1140 | 075 | 308 | 206 | 0,10,14 | - | 99 |
| 011 | 537 | -461 | 042 | - | - | 077 | - | 96 | 0,10,16 | - | -151 |
| 013 | 329 | 437 | 044 | 260 | 256 | 079 | - | -47 | 0,11,1 | 125 | 212 |
| 015 |  | -17 | 046 | - | - | 0,7,11 | 148 | - 144 | 0,11,3 | 125 | -57 |
| 017 | 271 | 164 | 048 | 424 | 313 | 0,7,13 | 279 | 259 | 0,11,5 | - | -76 |
| 019 | 244 | -206 | 0,4,10 |  |  | 0,7,15 | - | -20 | 0,11,7 | 127 | -119 |
| 0,1,11 |  | -16 | 0,4,12 | 268 | 184 | 0,7,17 | - | - | 0,11,9 | 120 | 123 |
| 0,1,13 | 168 | -219 | 0,4,14 | - |  | 0,7,19 |  | -115 | 0,11,11 |  | 49 |
| 0,1,15 | 121 | 181 | 0,4,16 | 275 | 321 | 080 | 283 | 219 | 0,11,13 | 65 | 37 |
| 0,1,17 | - | -116 | 0,4,18 |  |  | 082 |  |  | 0,11,15 |  | -98 |
| 0,1,19 | - | 81 | 0,4,20 | - | 37 | 084 | 529 | 557 | 0,12,0 | 223 | 232 |
| 0,1,21 | - | 3 | 051 | 542 | -568 | 086 |  | - | 0,12,2 |  |  |
| 020 | - | -55 | 053 | 135 | 102 | 088 | - | -133 | 0,12,4 | 208 | 150 |
| 022 | 646 | 584 | 055 | 310 | 334 | 0,8,10 | - |  | 0,12,6 |  |  |
| 024 | 481 | -377 | 057 | 432 | 312 | 0,8,12 | 353 | 332 | 0,12,8 | 56 | 38 |
| 026 | 540 | 483 | 059 | 445 | -319 | 0,8,14 |  |  | 0,12,10 |  |  |
| 028 | 621 | 514 | 0,5,11 | - | - 141 | 0,8,16 | - | 68 | 0,12,12 | 86 | 106 |
| 0,2,10 | 361 | 346 | 0,5,13 | - | -65 | 0,8,18 | - |  | 0,12,14 |  |  |
| 0,2,12 | 200 | -232 | 0,5,15 | 197 | 230 | 091 | 93 | -72 | 0,13,1 | - | 44 |
| 0,2,14 | 280 | 228 | 0,5,17 | 100 | -80 | 093 | 315 | 308 | 0,13,3 | 187 | 214 |
| 0,2,16 |  | 14 | 0,5,19 | - | 18 | 095 | - | - 101 | 0,13,5 | 127 | -119 |
| 0,2,18 | 133 | 136 | 0,5,21 | - | 37 | 097 | - | -23 | 0,13,7 | - | -73 |
| 0,2,20 | - | -46 | 060 | 604 | -721 | 099 | 174 | -6 | 0,13,9 | - | 54 |
| 0,2,22 | - | 75 | 062 | 368 | 391 | 0,9,11 | 114 | 77 | 0,13,11 | - | 85 |
| 031 | - | -187 | 064 | 343 | 387 | 0,9,13 | - | $-167$ | 0,14,0 | 53 | --38 |
| 033 | 770 | -959 | 066 | 484 | 340 | 0,9,15 | - | 40 | 0,14,2 | 92 | 95 |
| 035 | 645 | 485 | 068 | - | - 144 | 0,9,17 | - | -19 | 0,14,4 | - | -23 |
| 037 | 375 | 297 | 0,6,10 | 274 | 258 | 0,10,0 | 391 | -446 | 0,14,6 | 86 | 84 |
| 039 | - | -196 | 0,6,12 | 102 | 193 | 0,10,2 | 216 | 214 | 0,14,8 | - | 63 |
| 0,3,11 | 447 | -308 | 0,6,14 | 214 | 202 | 0,10,4 | 278 | 258 | 0,15,1 | 一 | 72 |
| 0,3,13 | 442 | 463 | 0,6,16 | 142 | -217 |  |  |  | 0,15,3 | - | -38 |

Structure factors are calculated with the axis origin at $2 / \mathrm{m}$.
by electron-density projections $\varrho_{0}(x, y)$ and $\varrho_{0}(y, z)$ (Fig. 1, 2). A $\varrho_{0}-\varrho_{P}(x, y)$ projection ( $\varrho_{P}=$ calculated electron-density projection due to the heavy atoms) showed that water molecules were in the positions of high density. However, their localization was not possible because the density values of the corresponding peaks were similar in magnitude to those of peaks ascribable to absorption and other experimental errors. The standard deviation of $\varrho_{0}(x, y)$ was $\sigma(\varrho)=3 \cdot 5 \mathrm{e} . \AA^{-2}$.
In $\varrho_{0}(x, y)$ the electron-density projection maxima corresponding to $\mathrm{Pb}, \mathrm{Sr}, \mathrm{I}_{1}, \mathrm{I}_{11}$ appear clearly; $\sigma(\varrho)$ was 6.29 e. $\AA^{-2}$.
Observed and calculated structure factors for $h k 0$ reflections are reported in Table 1. The temperature factor was found graphically to be $B=2.7 \AA^{2}$ and the discrepancy index was $R_{h k 0}=0.08$. (Unobserved reflections are not included.) In Table 2 are reported the observed and calculated structure factors for $0 k l$ reflections. The temperature factor was $B=3.5 \AA^{2}$. The value of $R_{0 k l} 0 \cdot 20$, was higher than that for $h k 0$ reflections. Nevertheless it did not seem worthwhile to try to reduce the value. In fact, though the range of $R$ was greatly changed by introducing absorption corrections derived from fairly accurate graphical methods, the plot of $R_{h k 0}$ versus $x_{\mathrm{IIII}}$ and of $R_{0 k l}$ versus $z_{\mathrm{I}_{\mathrm{I}}}$ always showed a minimum at nearly the same value of $x_{\mathrm{III}}$ and $z_{\mathrm{I}_{\mathrm{I}}}$ respectively. Therefore these minima are independent of absorption and depend only on $x_{\mathrm{I}_{\text {II }}}$ and $z_{\mathrm{I}_{\mathrm{I}}}$. Intensity data very roughly corrected never changed the $x_{\mathrm{I}_{\text {II }}}$ coordinate by more than 0.004 . Moreover $\varrho_{0}-\varrho_{P}(x, y)$ shows that experimental errors in the $F_{o}$ 's are always too high for the determination of the water coordinates, even if $R$ is low. The index $R$ is consequently an unreliable indication of the accuracy. This implies that we cannot use any expression depending on $F_{o}-F_{c}$ to assess the standard deviation in atomic coordinates. Therefore we assumed the value $\pm 0.004$ as standard deviation for atomic coordinates of both iodine atoms. The final coordinates are reported in Table 3.

Table 3. Atomic coordinates

|  | Number <br> of <br> atoms | Crys- <br> tall. <br> position | $x / a$ | $y / a$ | $z / c$ |
| :--- | :---: | :---: | :---: | :---: | :---: |
| Atom | 4 | $a$ | 0 | 0 | 0 |
| Sr | 8 | $d$ | 0 | 0.250 | 0.625 |
| Pb | 8 | $e$ | 0 | 0 | 0.310 |
| I | 16 | $g$ | 0.309 | 0.309 | 0 |
| I | 8 | $e$ | 0 | 0 | 0.141 |
| $\mathrm{H}_{2} \mathrm{O}_{\mathrm{I}}$ | 16 | $h$ | 0.208 | 0.000 | 0 |
| $\mathrm{H}_{2} \mathrm{O}_{\text {II }}$ | 16 | $b$ | 0 | 0 | 0.500 |
| $\mathrm{H}_{2} \mathrm{O}_{\text {III }}$ | 4 |  |  |  |  |

Axis origin at $\overline{4} 2 m$. Add 0 to $x / a ; 0.250$ to $y / a ;-0.125$ to $z / c$ for the origin at $2 / m$.

The numbers in italics are the variable independent parameters.

## Discussion of the structure

The structure of $\mathrm{SrPb}_{2} \mathrm{I}_{6} .7 \mathrm{H}_{2} \mathrm{O}$ is shown diagrammatically in Fig. 3. Distances between Pb and I


Fig. 3. Structure of $\mathrm{SrPb}_{2} \mathrm{I}_{6} .7 \mathrm{H}_{2} \mathrm{O}$. A quarter of a unit cell is represented.
$\left(\mathrm{Pb}-\mathrm{I}_{\text {II }}=3 \cdot 20 \pm 0 \cdot 02, \mathrm{~Pb}-\mathrm{I}_{\mathrm{I}}=3 \cdot 14 \pm 0 \cdot 04 \AA\right.$ ) are somewhat shorter than the sum of the ionic radii (Pauling, 1927). The covalent bond length of $\mathrm{Pb}-\mathrm{I}$ was found by Lister \& Sutton (1941) to be $2 \cdot 79 \pm 0.02 \AA$ in gaseous $\mathrm{PbI}_{2}$. Our value thus favours a partially covalent character of $\mathrm{Pb}-\mathrm{I}$ bonds, in agreement with the considerations of Powell \& Tasker (1937) concerning the $\mathrm{Pb}-\mathrm{Br}$ bonds in $\mathrm{NH}_{4} \mathrm{~Pb}_{2} \mathrm{Br}_{5}$; in this compound, however, coordination around Pb is not octahedral.
The shortening of the $\mathrm{Pb}-\mathrm{I}$ bond length as compared with the sum of the ionic radii is relatively small and is probably due to the sharing of the iodine atoms by two lead atoms. Actually it was observed in $\mathrm{NbCl}_{5}$ (Zalkin \& Sands, 1958) and in $\mathrm{Cs}_{3} \mathrm{Cr}_{2} \mathrm{Cl}_{9}$ (Wessel \& IJdo, 1957) that chlorine atoms shared by two octahedra are farther from the metal than the unshared ones $(2 \cdot 56 \AA$ for shared Cl against $2 \cdot 25-2 \cdot 30 \AA$ for unshared Cl in $\mathrm{NbCl}_{5}$ or $2.52 \AA$ for shared Cl against
$2 \cdot 34 \AA$ for unshared Cl in $\mathrm{Cs}_{3} \mathrm{Cr}_{2} \mathrm{Cl}_{9}$ ). The shortening of the metal-halogen bond length was observed in 1927 (Ferrari \& Baroni) in compounds of divalent transition metals containing chlorine atoms forming octahedral complexes with metals. It was confirmed by later researches (Brasseur \& Pauling, 1938; MacGillavry et al., 1939).

The main part of the structure are the octahedra $\left[\mathrm{PbI}_{6 / 2}\right]$ joined to one another through two opposite corners (occupied by $I_{I}$ ) to form linear chains throughout the whole crystal. There are two sets of such chains; the set 1 (see Fig. 4) includes chains parallel to [010] and the set 2 (see Fig. 4) includes chains parallel to [100]. The axes [010] and [100] can be interchanged. Chains in the same set are shifted $\frac{1}{2} a, \frac{1}{2} a, \frac{1}{2} c$. The two sets are joined to one another through the remaining corners of the octahedra (occupied by $\mathrm{I}_{\text {II }}$ ). It follows that it is made up of a


Fig. 4. Part of the three-dimensional framework formed of [ $\mathrm{PbI}_{6 / 2}$ ] octahedra sharing all corners. Corners not labelled are occupied by $\mathrm{I}_{\mathrm{II}}$. The arrows from $\mathrm{I}_{1}$ show the next bond $\mathrm{I}_{\mathrm{I}}-\mathrm{Pb}$ in the chain. The numbers 1,2 indicate the chains paraliel to [010] and [100] respectively.
three-dimensional framework formed of octahedra which share all corners (Fig. 4).

Each iodine atom is bound to two lead atoms. The angles between the two bond directions are practically equal in $\mathrm{I}_{\mathrm{I}}$ and in $\mathrm{I}_{\mathrm{II}}$ : $<\mathrm{Pb}-\mathrm{I}_{\mathrm{I}}-\mathrm{Pb}=139^{\circ} \pm 2^{\circ}$ : $<\mathrm{Pb}-\mathrm{I}_{\mathrm{II}}-\mathrm{Pb}=140^{\circ} \pm 2^{\circ}$. Distances $\mathrm{I}-\mathrm{I}$, which are the edges of the octahedra, are: $\mathrm{I}_{\mathrm{II}} \mathrm{I}_{\mathrm{II}}=4.55 \pm 0.04 \AA$; $\mathrm{I}_{\mathrm{I}}-\mathrm{I}_{\mathrm{II}}=4 \cdot 40 \pm 0.03 \AA ; \mathrm{I}_{\mathrm{I}^{\prime}}-\mathrm{I}_{\Pi^{\prime \prime}}=4.50 \pm 0.04 \AA ; \mathrm{I}_{\mathrm{I}^{\prime}}-\mathrm{I}_{\Pi}=$ $4.57 \pm 0.03 \AA$.

The empty spaces of the framework formed by the octahedra $\left[\mathrm{PbI}_{6 / 2}\right]$ are occupied by the strontium ions and by the water molecules. The coordinates of the water molecules could be determined only approximately, due to the presence of the heavy atoms and to absorption errors. Their octahedral arrangement around each strontium ion agrees well with the packing of atoms. The octahedron $\mathrm{Sr}\left(\mathrm{OH}_{2}\right)_{6}$ is of the same type as is found in $\mathrm{SrI}_{2} .6 \mathrm{H}_{2} \mathrm{O}$ and other related compounds (Hermann, 1931): assuming a distance $\mathrm{Sr}-\mathrm{OH}_{2}$ $=2.45 \AA$ we find $\mathrm{H}_{2} \mathrm{O}_{\mathrm{I}}-\mathrm{I}_{\mathrm{I}}=2.93 \AA$.

The distance between $\mathrm{H}_{2} \mathrm{O}_{\text {III }}$ and its neighbouring atoms are: $\mathrm{H}_{2} \mathrm{O}_{\mathrm{III}}-\mathrm{I}_{\mathrm{I}}=3 \cdot 38 \pm 0.02 \AA$ and $\mathrm{H}_{2} \mathrm{O}_{\mathrm{III}}-\mathrm{Pb}=$ $3 \cdot 65 \pm 0.02 \AA$. These distances exclude chemical bonds. Hence $\mathrm{H}_{2} \mathrm{O}_{\text {III }}$ is to be considered as of zeolyte- or clathrate-type.

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